

## On Infinite Series of Hypergeometric Function of Three Variables

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### Abstract

This paper deals with an integral transformation involving Whittaker function  $M_{k,m}$  into Lauricella's function  $F_2$ . Particular cases of this transformation correspond to a general triple hypergeometric function  $F^{(3)}$  and several known and new transformations are also discussed.

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### 1. Introduction and Definition

A Whittaker function  $M_{k,m}$  was introduced by Whittaker [5] (see also Whittaker and Watson [6]) in terms of Confluent hypergeometric function  ${}_1F_1$  (or Kummer's function)

$$M_{k,m}(x) = x^{m+1/2} e^{(-1/2)x} {}_1F_1\left(\frac{1}{2} + m - k, 2m + 1; x\right) \quad (1.1)$$

and the Horn's function  $H_4$  [6;p.39] is defined as

$$H_4[\alpha, \beta, \gamma, \delta; x, y] = \sum_{m,n=0}^{\infty} \frac{(\alpha)_{2m+n} (\beta)_n x^m y^n}{(\gamma)_m (\delta)_n m! n!} \quad (1.2)$$

where  $(a)_n = \frac{\Gamma(a+n)}{\Gamma(a)}$ ,  $|x| < r$ ,  $|y| < s$  and  $4r = (s-1)^2$ .

A general triple hypergeometric series  $F^{(3)}[x, y, z]$  of Srivastava [7; p.69(39)] defined as

$$F^{(3)} \left[ \begin{array}{l} (a) \quad :: (b); (b'); (b''); (c); (c'); (c'') \quad ; \\ (e) \quad :: (g); (g'); (g''); (h); (h'); (h'') \quad ; \end{array} \quad \begin{array}{l} x, y, z \\ \end{array} \right]$$

$$= \sum_{m,n,p=0}^{\infty} \frac{((a))_{m+n+p}((b))_{m+n}((b'))_{n+p}((b''))_{p+m}((c))_m((c'))_n((c''))_p}{((e))_{m+n+p}((g))_{m+n}((g''))_{n+p}((g''))_{p+m}((h))_m((h'))_n((h''))_p} \frac{x^m y^n z^p}{m! n! p!}, \quad (1.3)$$

with, as usual,  $(\alpha)$  abbreviates the array of  $A$ -parameters

$$\alpha_1, \alpha_2 \cdots \alpha_A, \quad ((\alpha))_m = \prod_{j=1}^A (\alpha_j)_m = \prod_{j=1}^A \frac{\Gamma(\alpha_j + m)}{\Gamma(\alpha_j)}. \quad (1.4)$$

## 2. Integral Transform

We establish the following integral

$$\int_0^{\infty} t^{\sigma-1/2} e^{-pt} {}_1F_2 \left( \begin{array}{l} \eta \quad ; \\ \gamma, \delta \quad ; \end{array} \quad -x^2 t^2 \right) M_{\kappa_1, \mu_1 - 1/2}(\beta_1 t) M_{\kappa_2, \mu_2 - 1/2}(\beta_2 t) dt$$

$$= \sum_{r=0}^{\infty} \frac{(\eta)_r (-1)^r x^{2r} \beta_1^{\mu_1} \beta_2^{\mu_2} \Gamma(a + 2r)}{(\gamma)_r (\delta)_r r! (p + \frac{\beta_1}{2} + \frac{\beta_2}{2})^{a+2r}}$$

$$F_2 \left[ a + 2r, \mu_1 - k_1, \mu_2 - k_2, 2\mu_1, 2\mu_2; \frac{\beta_1}{p + \frac{\beta_1}{2} + \frac{\beta_2}{2}}, \frac{\beta_2}{p + \frac{\beta_1}{2} + \frac{\beta_2}{2}} \right] \quad (2.1)$$

where  $a = \sigma + \mu_1 + \mu_2 + \frac{1}{2}$ ,  $Re(p + \frac{\beta_1}{2} + \frac{\beta_2}{2}) > 0$  and  $Re(a) > 0$  and Appell's function  $F_2$  [7;p.53] is defined as

$$F_2 [a, b, b'; c, c'; x, y] = \sum_{m,n=0}^{\infty} \frac{(a)_{m+n} (b)_m (b')_n x^m y^n}{(c)_m (c')_n m! n!}.$$

Equation (2.1) can be established by expanding  ${}_1F_2$  in series and integrating term by term with the help of a result [3;p.216(14)].

Setting  $k_1 = k_2 = 0$  in (2.1), using a result of Bailey [10;p.11] involving  $F_2$  and  $F_4$  [2;p.224]

$$F_2 \left[ \alpha, \beta - \frac{1}{2}, \beta' - \frac{1}{2}, 2\beta - 1, 2\beta' - 1; 2x, 2y \right]$$

$$= (1 - x - y)^{-\alpha} F_4 \left[ \frac{1}{2}\alpha, \frac{1}{2}\alpha + \frac{1}{2}, \beta, \beta'; \frac{x^2}{(1 - x - y)^2}, \frac{y^2}{(1 - x - y)^2} \right] \quad (2.2)$$

Expanding  $F_4$  in series and making an appeal of Legendre's duplication formula [2;p.5(15)]

$$\left(\frac{1}{2}a\right)_m \left(\frac{1}{2}a + \frac{1}{2}\right)_m = 2^{-2m} (a)_{2m}, m = 0, 1, 2, \dots$$

in conjugation with (1.3), we get

$$\int_0^\infty t^{\sigma-1/2} e^{-pt} {}_1F_2 \left( \begin{matrix} \eta & ; \\ \gamma, \delta & ; \end{matrix} -x^2 t^2 \right) M_{0, \mu_1-1/2}(\beta_1 t) M_{0, \mu_2-1/2}(\beta_2 t) dt$$

$$= \frac{\beta_1^{\mu_1} \beta_2^{\mu_2} \Gamma(a)}{p^a} F^{(3)} \left[ \begin{matrix} \frac{1}{2}a, \frac{1}{2}a + \frac{1}{2} & :: \text{---} : \eta; \text{---}; \text{---}; \text{---}; \\ \text{---} & :: \text{---} : \gamma, \delta; \mu_1 + \frac{1}{2}; \mu_2 + \frac{1}{2}; \end{matrix} \frac{-4x^2}{p^2}, \frac{\beta_1^2}{4p^2}, \frac{\beta_2^2}{4p^2} \right] \quad (2.3)$$

where  $a = \sigma + \mu_1 + \mu_2 + \frac{1}{2}$ ,  $Re(a) > 0$  and  $Re\left(p + \frac{\beta_1}{2} + \frac{\beta_2}{2}\right) > 0$ .

Comparing of (2.3) to (2.1) with  $k_1 = k_2 = 0$  would yield an expansion

$$\sum_{r=0}^\infty \frac{(\eta)_r (-1)^r (\alpha)_{2r}}{(\gamma)_r (\delta)_r r!} F_2 [\alpha + 2r, \sigma, \omega, 2\sigma, 2\omega; y, z]$$

$$= \left(1 - \frac{1}{2}y - \frac{1}{2}z\right)^{-\alpha} F^{(3)} \left[ \begin{matrix} \frac{1}{2}a, \frac{1}{2}a + \frac{1}{2} & :: \text{---} : \eta; \text{---}; \text{---}; \text{---}; \\ \text{---} & :: \text{---} : \gamma, \delta; \sigma + \frac{1}{2}; \omega + \frac{1}{2}; \end{matrix} \frac{-16x^2}{(2-y-z)^2}, \frac{y^2}{(2-y-z)^2}, \frac{z^2}{(2-y-z)^2} \right] \quad (2.4)$$

In view of a known transformation [1;p.381].

$$F_2 [\alpha, \beta, \beta', 2\beta, 2\beta'; 2x, y] = (1 - x)^{-\alpha} H_4 \left[ \alpha, \beta', \beta + \frac{1}{2}, 2\beta'; \frac{x^2}{4(1 - x)^2}, \frac{y}{1 - x} \right], \quad (2.5)$$

equation (2.4) gives an infinite sum of  $H_4$  in the form

$$\sum_{r=0}^\infty \frac{(\eta)_r (-1)^r (\alpha)_{2r}}{(\gamma)_r (\delta)_r r!} \left(1 - \frac{y}{2}\right)^{-2r} H_4 \left[ \alpha + 2r, \omega, \sigma + \frac{1}{2}, 2\omega; \frac{y^2}{4(2 - y)^2}, \frac{2z}{2 - y} \right]$$

$$= \left( \frac{2-y-z}{2-y} \right)^{-\alpha} F^{(3)} \left[ \begin{array}{c} \frac{1}{2}a, \frac{1}{2}a + \frac{1}{2} \\ \hline \end{array} \begin{array}{c} \text{---} : \eta; \text{---}; \text{---}; \text{---}; \\ \frac{-16x^2}{(2-y-z)^2}, \frac{y^2}{(2-y-z)^2}, \\ \text{---} : \gamma, \delta; \sigma + \frac{1}{2}; \omega + \frac{1}{2}; \\ \frac{z^2}{(2-y-z)^2} \end{array} \right] \quad (2.6)$$

For  $z \rightarrow 0$ , equation (2.4) gives

$$\sum_{r=0}^{\infty} \frac{(\eta)_r (-1)^r (\alpha)_{2r}}{(\gamma)_r (\delta)_r r!} {}_2F_1 \left( \begin{array}{c} \alpha + 2r, \sigma \\ 2\omega \end{array} ; y \right) x^{2r} \\ = \left( 1 - \frac{1}{2}y \right)^{-\alpha} F_{0;2;1}^{2:1;0} \left[ \begin{array}{c} \frac{1}{2}a, \frac{1}{2}a + \frac{1}{2} \\ \hline \end{array} \begin{array}{c} \text{---} : \eta; \text{---}; \text{---}; \text{---}; \\ \frac{-16x^2}{(2-y)^2}, \frac{y^2}{(2-y)^2}, \\ \text{---} : \gamma, \delta; \sigma + \frac{1}{2}; \text{---}; \end{array} \right] \quad (2.7)$$

$$\sum_{r,k=0}^{\infty} \frac{(\eta)_r (-1)^r (\alpha)_{2r+k} (\sigma)_k y^k x^{2r}}{(\gamma)_r (\delta)_r (2\omega)_k r! k!} \\ = \left( 1 - \frac{1}{2}y \right)^{-\alpha} F_{0;2;1}^{2:1;0} \left[ \begin{array}{c} \frac{1}{2}a, \frac{1}{2}a + \frac{1}{2} \\ \hline \end{array} \begin{array}{c} \text{---} : \eta; \text{---}; \text{---}; \text{---}; \\ \frac{-16x^2}{(2-y)^2}, \frac{y^2}{(2-y)^2}, \\ \text{---} : \gamma, \delta; \sigma + \frac{1}{2}; \text{---}; \end{array} \right] \quad (2.8)$$

On setting  $\eta = \gamma$ , the result (2.8) reduces to known result Pathan [9;p.41].

On setting  $k_2 = -\mu_2$  in (2.1), the result reduces to

$$\beta_2^{\mu_2} \int_0^{\infty} t^{\sigma+\mu_2-1/2} e^{-(p-1/2\beta_2)t} {}_1F_2 \left( \begin{array}{c} \eta \\ \gamma, \delta \end{array} ; -x^2 t^2 \right) M_{\kappa_1, \mu_1-1/2}(\beta_1 t) dt \\ = \beta_1^{\mu_1} \beta_2^{\mu_2} \sum_{r=0}^{\infty} \frac{(\eta)_r (-1)^r x^{2r} \Gamma(a+2r)}{(\gamma)_r (\delta)_r r! (p + \frac{\beta_1}{2} - \frac{\beta_2}{2})^{a+2r}} {}_2F_1 \left[ a+2r, \mu_1 - k_1, 2\mu_1; \frac{\beta_1}{p + \frac{\beta_1}{2} - \frac{\beta_2}{2}} \right]$$

$$= \frac{\beta_1^{\mu_1} \beta_2^{\mu_2} \Gamma(a)}{(p + \frac{\beta_1}{2} - \frac{\beta_2}{2})^a} \sum_{r,s=0}^{\infty} \frac{(\eta)_r (-1)^r (a)_{2r+s} (\mu_1 - k_1)_s x^{2r} y^s}{(\gamma)_r (\delta)_r (p + \frac{\beta_1}{2} - \frac{\beta_2}{2})^{2r} r! s!} \tag{2.9}$$

On setting  $\eta = \gamma, \delta = \mu, x \rightarrow \frac{x}{2}$  and  $x = 2\alpha$  in (2.9), the result reduces to known result Pathan [9;p.42].

### 3. Series Expansion

The generalized hypergeometric series of a power of t is given by [4;p.417].

$$\sum_{k=0}^{\infty} \frac{t^k}{k!} {}_pF_q \left( \begin{matrix} -k, (a_p) \\ (b_q) \end{matrix} ; x \right) = e^t {}_pF_q \left( \begin{matrix} (a_p) \\ (b_q) \end{matrix} ; -tx \right) \tag{3.1}$$

For  $p = 1, q = 2$ , equation (3.1) reduces to

$$\sum_{k=0}^{\infty} \frac{t^k}{k!} {}_2F_2 \left( \begin{matrix} -k, (a_1) \\ b_1, b_2 \end{matrix} ; x \right) = e^t {}_2F_2 \left( \begin{matrix} a_1 \\ b_1, b_2 \end{matrix} ; -tx \right) \tag{3.2}$$

Multiplying both the sides of (3.2)  $t^{\sigma-1/2} e^{-pt} M_{0,\mu_1-1/2}(\beta_1 t) M_{0,\mu_2-1/2}(\beta_2 t)$  and integrating with respect to t from zero to infinity with the help of the result (2.3) and [3;p.216(14)], we obtain

$$\begin{aligned} & \sum_{k=0}^{\infty} \frac{(a)_k}{p^k k!} F^{(3)} \left[ \begin{matrix} - & :: & \frac{1}{2}(a+k), \frac{1}{2}(a+k) + \frac{1}{2}; - & : & -k, a_1; \_ & ; \_ & ; \\ & & & & & & & & x, \frac{\beta_1^2}{4p^2}, \frac{\beta_2^2}{4p^2} \end{matrix} \right] \\ &= \left( \frac{p}{p-1} \right)^a \sum_{m=0}^{\infty} \frac{(-1)^m (a_1)_m (a)_m x^m}{(b_1)_m (b_2)_m (p-1)^m m!} F^{(2)} \\ & \times \left[ \begin{matrix} \frac{1}{2}(a+m), \frac{1}{2}(a+m) + \frac{1}{2} : \_ & ; \_ & ; & \frac{\beta_1^2}{4(p-1)^2}, \frac{\beta_2^2}{4(p-1)^2} \\ \_ & : & \mu_1 + \frac{1}{2}; \mu_2 + \frac{1}{2}; \end{matrix} \right] \tag{3.3} \end{aligned}$$

On setting  $x = 0$  in (3.3), the result reduces to

$$\sum_{k=0}^{\infty} \frac{(a)_k}{p^k k!} F^{(2)} \left[ \begin{matrix} - & :: & \frac{1}{2}(a+k), \frac{1}{2}(a+k) + \frac{1}{2}; - & : & -; -; \_ & ; \_ & ; \\ & & & & & & & & \frac{\beta_1^2}{4p^2}, \frac{\beta_2^2}{4p^2} \end{matrix} \right]$$

$$= \left( \frac{p}{p-1} \right)^a F^{(2)} \left[ \begin{array}{c} \frac{1}{2}(a+m), \frac{1}{2}(a+m) + \frac{1}{2} : \text{---}; \text{---}; \\ \frac{\beta_1^2}{4(p-1)^2}, \frac{\beta_2^2}{4(p-1)^2} \\ \text{---} : \mu_1 + \frac{1}{2}; \mu_2 + \frac{1}{2}; \end{array} \right] \quad (3.4)$$

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