

The Quenching Behavior of a Semilinear Parabolic Equation with a Singular Boundary Outflux *

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Abstract

In this paper, we consider the quenching behavior of the solution of a semilinear parabolic equation with a singular boundary outflux. First we prove a finite-time quenching for the solution. then, we assert that quenching occurs on the boundary under certain conditions .Futhermore,we show that the time derivative blows up at a quenching point. Finally, we get a quenching rate and a lower bound for the quenching time.

Keywords: semilinear parabolic equation, singular boundary outflux, quenching, maximum principles

1. INTRODUCTION

In this paper, we consider the quenching behavior of solutions of the following semilinear parabolic equation with a singular boundary outflux:

$$\begin{cases} u_t = u_{xx} + f(x)(1-u)^{-p}, & 0 < x < 1, 0 < t < T, \\ u_x(0, t) = 0, u_x(1, t) = -u^{-q}(1, t), & 0 < t < T, \\ u(x, 0) = u_0(x), & 0 \leq x \leq 1, \end{cases} \quad (1.1)$$

where p, q are positive constants and $T \leq \infty$. The initial function $u_0 : [0, 1] \rightarrow (0, 1)$

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satisfies the compatibility conditions $u_0'(0) = 0, u_0'(1) = -u_0^{-q}(1)$. , And $f(x)$ is a continuous and positive function. Throughout this paper, we assume that the initial function u_0 satisfies inequalities

$$u_0''(x) + f(x)(1 - u_0(x))^{-p} \geq 0. \quad (1.2)$$

$$u_0'(x) \leq 0. \quad (1.3)$$

$$f'(x) \leq 0. \quad (1.4)$$

Since 1975, quenching problems with various boundary conditions have been studied extensively (cf. the survey by Chan [1, 2] and Kirk and Roberts [14] and [3, 4, 6–9, 12]). In the literature, quenching problems have been less studied with two nonlinear heat sources. We give as examples five of these papers.

Chan and Yuen [5] considered the problem

$$\begin{cases} u_t = u_{xx}, & \text{in } \Omega, \\ u_x(0, t) = (1 - u(0, t))^{-p}, u_x(a, t) = (1 - u(a, t))^{-q}, & 0 < t < T, \\ u(x, 0) = u_0(x), 0 \leq u_0(x) < 1, & \text{in } \Omega, \end{cases}$$

where $a, p, q > 0, T \leq \infty, D = (0, a), \Omega = D \times (0, T)$.

They showed that $x = a$ is the unique quenching point in finite time if u_0 is a lower solution, and u_t blow up at quenching. Finally, they obtained criteria for nonquenching and quenching by using the positive steady states.

Burhan selcuk and Nuri ozalp [15] considered the problem

$$\begin{cases} u_t = u_{xx} + (1 - u)^{-p}, & 0 < x < 1, 0 < t < T, \\ u_x(0, t) = 0, u_x(1, t) = -u^{-q}(1, t), & 0 < t < T, \\ u(x, 0) = u_0(x), & 0 \leq x \leq 1, \end{cases}$$

where $p, q > 0, T \leq \infty$.

They showed $x = 0$ is the unique quenching point in finite time if u_0 satisfies $u_{xx}(x, 0) + (1 - u(x, 0))^{-p} \geq 0, u_x(x, 0) \leq 0$, and u_t blow up at quenching time. Finally, they obtained a lower bound for the quenching time.

Ozalp and Selcuk [11] considered the problem

$$\begin{cases} u_t = u_{xx} + (1 - u)^{-p}, & 0 < x < 1, 0 < t < T, \\ u_x(0, t) = 0, u_x(1, t) = (1 - u(1, t))^{-p}, & 0 < t < T, \\ u(x, 0) = u_0(x), & 0 \leq x \leq 1, \end{cases}$$

where $p, q > 0, T \leq \infty$.

They showed $x = 1$ is the unique quenching point in finite time if u_0 satisfies $u_{xx}(x, 0) + (1 - u(x, 0))^{-p} \geq 0, u_x(x, 0) \geq 0$, and u_t blow up at quenching time. Finally, they obtained a quenching rate and a lower bound for the quenching time.

Lin and Wang [13] considered the problem

$$\begin{cases} u_t = u_{xx} + (1 - u)^{-p}, & 0 < x < 1, 0 < t < T, \\ u_x(0, t) = 0, u_x(1, t) = -u^{-q}(1, t), & 0 < t < T, \\ u(x, 0) = u_0(x), & 0 \leq x \leq 1, \end{cases}$$

where $p, q > 0, T \leq \infty$.

they showed that the solution have a finite time blow-up and obtained the exact blow-up rates for the necessary and sufficient conditions. They also proved that the blow-up will occur only at the boundary $x = 1$. Then, by applying the well-known method of Giga-Kohn, they derived the time asymptotic of solutions near the blow-up time. Finally, they proved that the blow-up was complete.

Yuanhong Zhi [16] consider the problem

$$\begin{cases} u_t = u_{xx} + f(x)(1 - u)^{-p}, & 0 < x < 1, 0 < t < T, \\ u_x(0, t) = u^{-q}(0, t), u_x(1, t) = 0, & 0 < t < T, \\ u(x, 0) = u_0(x), & 0 \leq x \leq 1, \end{cases}$$

where $p, q > 0, T \leq \infty, f(x)$ is a non-negative function.

They showed $x = 0$ is the unique quenching point in finite time if u_0 satisfies $u_{xx}(x, 0) + f(x)(1 - u(x, 0))^{-p} \leq 0, u_x(x, 0) \geq 0, f'(x) \geq 0$, further, they obtained lower and upper bounds for the quenching time.

In this paper, a quenching problem with two types of singularity terms, namely, a source term $(1 - u)^{-p}$, the boundary outflux term $-u^{-q}$, and the effect of $f(x)$, are considered. In Section 2, we first show that quenching occurs in finite time under condition (1.2), Then we show that the only quenching point $x = 0$ under condition (1.2) and (1.3), Further we show that u_t blow up at the quenching point $x = 0$, In Section 3, we get a quenching rate and a lower bound for quenching time.

2. QUENCHING ON THE BOUNDARY AND BLOW-UP OF U_T

Definition 1 A solution $u(x,t)$ of problem (1.1) is said to quench if there exists a finite time T such that

$$\lim_{t \rightarrow T^-} \max\{u(x,t) : 0 \leq x \leq 1\} \rightarrow 1 \text{ or } \lim_{t \rightarrow T^-} \min\{u(x,t) : 0 \leq x \leq 1\} \rightarrow 0$$

From now on, we denote the quenching time of problem (1.1) with T .

The main results of this paper read as follows.

lemma 1 if u_0 satisfies (1.3), then $u_x < 0$ in $[0, 1] \times [0, T)$.

proof let $w = u_x$, then w satisfies

$$\begin{cases} w_t = w_{xx} + f'(x)(1-u)^{-p} + f(x)p(1-u)^{-p-1}w, & 0 < x < 1, 0 < t < T, \\ w(0,t) = 0, w(1,t) = -u^{-q}(1,t), & 0 < t < T, \\ w(x,0) = u'_0(x) & 0 \leq x \leq 1, \end{cases}$$

since $u'_0(x) \leq 0$, then $w(x,0) \leq 0$. besides, $w(0,t) = 0$, $w(1,t) = -u^{-q}(1,t) < 0$, thus by the maximum principle $u_x(x,t) < 0$ in $(0, 1] \times (0, T)$

lemma 2 If u_0 satisfies (1.2), then $u_t \geq 0$ in $[0, 1] \times [0, T)$.

proof we give the proof by utilizing Lemma 3.1 in [10]. Let $v = u_t(x,t)$. Then $v(x,t)$ satisfies

$$\begin{cases} v_t = v_{xx} + f(x)p(1-u)^{-p-1}v, & 0 < x < 1, 0 < t < T, \\ v_x(0,t) = 0, v_x(1,t) = qu^{-q-1}(1,t)v(1,t), & 0 < t < T, \\ v(x,0) = u_{xx}(x,0) + f(x)(1-u(x,0))^{-p} \geq 0, & 0 \leq x \leq 1, \end{cases}$$

For any fixed $\tau \in (0, T)$, let

$$L = \max_{0 \leq x \leq 1, 0 \leq t \leq \tau} \left(\frac{1}{2} qu^{-q-1}(x,t) \right), M = 2L + 4L^2 + \max_{0 \leq x \leq 1, 0 \leq t \leq \tau} (f(x)p(1-u(x,t))^{-p-1}). \quad (2.1)$$

Set $\omega(x,t) = e^{-Mt-Lx^2}v(x,t)$. Then ω satisfies

$$\begin{cases} \omega_t = \omega_{xx} + 4Lx\omega_x + c\omega, & 0 < x < 1, 0 < t < T, \\ \omega_x(0,t) = 0, \omega_x(1,t) = d(t)(1,t)\omega(1,t), & 0 < t < T, \\ \omega(x,0) \geq 0, & 0 \leq x \leq 1, \end{cases}$$

where

$$c = c(x,t) = 4L^2(x^2 - 1) + f(x)p(1-u(x,t))^{-p-1} - \max_{0 \leq x \leq 1, 0 \leq t \leq \tau} (f(x)p(1-u(x,t))^{-p-1}) \leq 0. \quad (2.2)$$

and

$$d(t) = - \max_{0 \leq x \leq 1, 0 \leq t \leq \tau} (qu^{-q-1}(x, t)) + qu^{-q-1}(1, t) \leq 0 \tag{2.3}$$

By the maximum principle and the Hopf lemma, we obtained that $\omega \geq 0$ in $[0, 1] \times [0, \tau]$. Thus $u_t(x, t) \geq 0$ in $[0, 1] \times [0, T)$.

Theorem 1 If u_0 satisfies (1.2), then there exists a finite time T such that the solution u of problem (1.1) quenches at time T .

proof We prove this result by contradiction. Assume on the contrary u can not quench at all time. Set

$$\gamma := -u^{-q}(1, 0) + \int_0^1 f(x)(1 - u(x, 0))^{-p} dx \geq 0, \tag{2.4}$$

Introduce a mass function $m(t) = \int_0^1 (1 - u(x, t)) dx, 0 < t < T$. Then

$$m'(t) = u^{-q}(1, 0) - \int_0^1 f(x)(1 - u(x, t))^{-p} dx \leq -\gamma. \tag{2.5}$$

by Lemma 2. Thus, $m(t) \leq m(0) - \gamma t$, which means that $m(T_0) = 0$ for some $T_0 (0 < T \leq T_0)$, which means that u quenches in a finite time.

Theorem 2 If u_0 satisfies (1.2)(1.3) and (1.4), then $x = 0$ is the only quenching point.

proof Define

$$K(x, t) = u_x + \epsilon(b_2 - x) \text{ in } [b_1, b_2] \times [\tau, T), \tag{2.6}$$

where $b_2 \in (0, 1], b_1 \in (0, b_2), \tau \in [0, T)$ and ϵ is a positive constant to be specified later. Then $K(x, t)$ satisfies

$$K_t - K_{xx} = f'(x)(1 - u)^{-p} + f(x)p(1 - u)^{-p-1}u_x. \tag{2.7}$$

since $u_x(x, t) < 0, f'(x) \leq 0$ in $(0, 1] \times [0, T)$. Thus, $K(x, t) - K_{xx} < 0$ in $(b_1, b_2) \times (\tau, T)$. Further, if ϵ is small enough,

$$K(b_1, t) = u_x(b_1, t) + \epsilon(b_2 - b_1) < 0, \tag{2.8}$$

$$K(b_2, t) = u_x(b_2, t) < 0, \quad (2.9)$$

for $t \in (\tau, T)$. By the maximum principle, we obtain that $K(x, t) < 0$ in $[b_1, b_2] \times [\tau, T]$, i.e., $u_x < -\epsilon(b_2 - x)$ for $(x, t) \in [b_1, b_2] \times [\tau, T]$. Integrating this with respect to x from b_1 to b_2 , we have

$$u(b_2, t) < u(b_1, t) - \frac{\epsilon(b_2 - b_1)^2}{2} < 1 - \frac{\epsilon(b_2 - b_1)^2}{2} < 1. \quad (2.10)$$

So u does not quench in $(0, 1]$. The theorem is proved.

Remark 1 Our hypothesis is reasonable, in other words, we can make a initial function to satisfies (1.2), (1.3) and the compatibility conditions, for example, let $u_0(x) = 0.9 - \frac{2}{3}x^{4.5}$, when $p = 9$ and $q = \log_{30/7}(3)$, $u_0(x)$ satisfies (1.2), (1.3) and the compatibility conditions

Theorem 3 if $p \geq 1$, then u_t blows up at the quenching point $x = 0$.

proof Suppose that u_t is bound on $[0, 1] \times [0, T)$, Then, there exists a positive constant M , such that $u_t < M$, That is

$$u_{xx} + f(x)(1 - u)^{-p} < M, \quad (2.11)$$

Since $f(x)$ is a continuous positive function, and $f'(x) \leq 0$, thus $f(x)$ is bounded. Then, there exists a positive constant a , such that $f(x) \geq a$, therefore the above inequality will be

$$u_{xx} + a(1 - u)^{-p} < M, \quad (2.12)$$

Multiplying this inequality by u_x , and integrating with respect to x from 0 to x . we have

$$a \ln(1 - u(0, t)) > -\frac{1}{2}u_x^2 + a \ln(1 - u(x, t)) + M[u(x, t) - u(0, t)], \quad (2.13)$$

for $p = 1$ and

$$a \frac{(1 - u(0, t))^{-p+1}}{-p+1} > -\frac{1}{2}u_x^2 + a \frac{(1 - u(0, t))^{-p+1}}{-p+1} + M[u(x, t) - u(0, t)] \quad (2.14)$$

for $p \neq 1$. We have, as $t \rightarrow T^-$ and $p \geq 1$, that the left-hand side tends to negative infinity, while the right-hand side is finite. This contradiction shows that u_t blows up at the quenching point $x = 0$.

3. A QUENCHING RATE AND A LOW BOUND FOR THE QUENCHING TIME

In this section, we get a quenching rate and a lower bound for the quenching time. Throughout this section, we assume that

$$u_x(x, 0) \leq -xu^{-q}(x, 0), \quad 0 \leq x \leq 1, \tag{3.1}$$

$$u_t(0, t) = u_{xx}(0, t) + f(0)(1 - u(0, t))^{-p}, \quad 0 < t < T. \tag{3.2}$$

Theorem 4 If u_0 satisfies (1.2), (1.3), (1.4), (3.1) and (3.2), then there exists a positive constant C_1 such that

$$u(0, t) \geq 1 - C_1(T - t)^{1/(p+1)}. \tag{3.3}$$

for t sufficiently close to T .

proof Define $J(x, t) = u_x + xu^{-q}$ in $[0, 1] \times [0, T]$. Then, $J(x, t)$ satisfies

$$J_t - J_{xx} = fp(1-u)^{-p-1}u_x + f'(1-u)^{-p} - fqx(1-u)^{-p}u^{-q-1} + 2qu^{-q-1}u_x - q(q+1)xu^{-q-2}u_x^2, \tag{3.4}$$

since $u_x < 0, f'(x) \leq 0, f(x) > 0, J_t - J_{xx} < 0$ in $(0, 1) \times [0, T]$. On the other hand, since

$$J(0, t) = 0, J(1, t) = 0, \tag{3.5}$$

for $t \in (0, T)$ and assumption (3.1), thus by the maximum principle, we obtained that $J(x, t) \leq 0$ for $(x, t) \in [0, 1] \times [0, T]$. Therefore

$$J_x(0, t) = \lim_{h \rightarrow 0^+} \frac{J(h, t) - J(0, t)}{h} = \lim_{h \rightarrow 0^+} \frac{J(h, t)}{h} \leq 0, \tag{3.6}$$

From (3.2), we get

$$J_x(0, t) = u_{xx}(0, t) + u^{-q}(0, t) = u_t(0, t) - f(0)(1 - u(0, t))^{-p} + u^{-q}(0, t) \leq 0, \tag{3.7}$$

and

$$u_t(0, t) \leq f(0)(1 - u(0, t))^{-p} \tag{3.8}$$

Integrating for t from t to T we get

$$u(0, t) \geq 1 - C_1(T - t)^{1/(p+1)}, \tag{3.9}$$

where $C_1 = [f(0)(p + 1)]^{1/p+1}$.

Remark 2 we can calculate a lower bound for the quenching time from Theorem 4, a lower bound is $(1 - u_0(0))^{p+1}/f(0)(p + 1)$ for quenching time T . If we choose, as in Remark 1, $u_0(x) = 0.9 - \frac{2}{3}x^{4.5}$, then we have $T = 10^{-11}$ for $p = 9$.

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